

Transient Thermoelastic Problem of Semi Infinite Rectangular Beam with Heat Generation

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Abstract- This paper is concerned with transient thermoelastic problem in which we need to determine the temperature distribution, displacement function and thermal stresses of a semi-infinite square beam when the boundary conditions are known. Integral transform techniques are used to obtain the solution of the problem.

Key Words: Semi-infinite rectangular beam, transient problem, Integral transform, heat source

I. INTRODUCTION

In 2003, Noda et al. [1] have published a book on Thermal Stresses, second edition. Dange et al. [2] discussed three dimensional inverse transient thermoelastic problem of a thin rectangular plate. Ghume et al. [3] studied deflection of a thick rectangular plate. Roy et al. [4] discussed transient thermoelastic problem of an infinite rectangular slab. Lamba, et al. [5] studied thermoelastic problem of a thin rectangular plate due to partially distributed heat supply. Sutar et al. [6] discussed an inverse thermoelastic problem of heat conduction with internal heat generation for the rectangular plate. Hiranwar et al. [7] studied thermal deflection of a thick clamped rectangular plate. Bagade et al. [8] discussed thermal stresses of a semi infinite rectangular beam. Jadhav et al. W [9] studied an inverse thermoelastic problem of a thin finite rectangular plate due to internal heat source. Singru et al. [10] discussed thermal stress analysis of a thin rectangular plate with internal heat source. Khobragade et al. [11] studied thermal stresses of a semi-infinite rectangular slab with internal heat generation.

In this paper, an attempt has been made to determine the temperature distribution, displacement function and thermal stresses at any point of a semi-infinite rectangular beam occupying the region $D : -a \leq x \leq a ; 0 \leq y \leq b, 0 \leq z \leq \infty$ with known boundary conditions. Here Marchi-Fasulo transforms and Fourier cosine transform techniques have been used to find the solution.

II. STATEMENT OF THE PROBLEM

Consider a semi-infinite rectangular beam occupying the space $D : -a \leq x \leq a ; 0 \leq y \leq b, 0 \leq z \leq \infty$. The displacement

components u_x, u_y and u_z in x and y and z directions respectively as Tanigawa et al. [1] are

$$u_x = \int_{-a}^a \left[\frac{1}{E} \left(\frac{\partial^2 U}{\partial y^2} + \frac{\partial^2 U}{\partial z^2} - \nu \frac{\partial^2 U}{\partial x^2} \right) + \lambda T \right] dx \quad (1)$$

$$u_y = \int_0^b \left[\frac{1}{E} \left(\frac{\partial^2 U}{\partial z^2} + \frac{\partial^2 U}{\partial x^2} - \nu \frac{\partial^2 U}{\partial y^2} \right) + \lambda T \right] dy \quad (2)$$

$$u_z = \int_0^\infty \left[\frac{1}{E} \left(\frac{\partial^2 U}{\partial x^2} + \frac{\partial^2 U}{\partial y^2} - \nu \frac{\partial^2 U}{\partial z^2} \right) + \lambda T \right] dz \quad (3)$$

where E, ν , and λ are the young's modulus, Poisson's ratio and the linear coefficient of the thermal expansion of the material of the beam respectively and $U(x, y, z, t)$ is the Airy's stress functions which satisfy the differential equation as Tanigawa et al. [1] is

$$\left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \right)^2 U(x, y, z, t) = -\lambda E \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \right) \times T(x, y, z, t) \quad (4)$$

where $T(x, y, z, t)$ denotes the temperature of a rectangular beam satisfy the following differential equation as Tanigawa et al. [1] is

$$\frac{\partial^2 T}{\partial x^2} + \frac{\partial^2 T}{\partial y^2} + \frac{\partial^2 T}{\partial z^2} + \frac{g(x, y, z, t)}{k} = \frac{1}{\alpha} \frac{\partial T}{\partial t} \quad (5)$$

where k is the thermal conductivity and α is the thermal diffusivity of the material, subject to initial condition

$$T(x, y, z, 0) = f(x, z, t) \quad (6)$$

The boundary conditions are

$$\left[T(x, y, z, t) + k_1 \frac{\partial T(x, y, z, t)}{\partial x} \right]_{x=a} = f_1(y, z, t) \quad (7)$$

$$\left[T(x, y, z, t) + k_2 \frac{\partial T(x, y, z, t)}{\partial x} \right]_{x=-a} = f_2(y, z, t) \quad (8)$$

$$[T(x, y, z, t)]_{y=0} = f_3(x, z, t) \tag{9}$$

$$[T(x, y, z, t)]_{y=b} = f_4(x, z, t) \tag{10}$$

$$\left. \frac{\partial T(x, y, z, t)}{\partial z} \right|_{z=0} = 0 \tag{11}$$

$$\left. \frac{\partial T(x, y, z, t)}{\partial z} \right|_{z=\infty} = h(x, y, t) \tag{12}$$

The stress components in terms of $U(x, y, z, t)$ Durge et al.[2] are given by

$$\sigma_{xx} = \left[\frac{\partial^2 U}{\partial y^2} + \frac{\partial^2 U}{\partial z^2} \right] \tag{13}$$

$$\sigma_{yy} = \left[\frac{\partial^2 U}{\partial z^2} + \frac{\partial^2 U}{\partial x^2} \right] \tag{14}$$

$$\sigma_{zz} = \left[\frac{\partial^2 U}{\partial x^2} + \frac{\partial^2 U}{\partial y^2} \right] \tag{15}$$

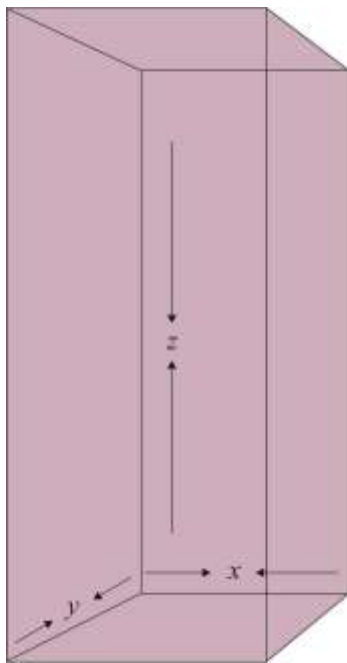


Figure 1: Geometry of the problem

The equations (1) to (15) constitute the mathematical formulation of the problem under consideration

III. SOLUTION OF THE PROBLEM

Applying finite Marchi-Fasulo transform defined in [2], finite Fourier sine transform and Fourier cosine transform to the equations, we get

$$\frac{d\bar{T}^*}{dt} + \alpha q^2 \bar{T}^* = \frac{\alpha g}{k} + \Psi$$

This is a linear equation whose solution is given by

$$\bar{T}^*(m, n, \eta, t) = e^{-\alpha q^2 t} \left(\bar{f}^* + \int_0^t \left[\frac{\alpha g}{k} + \Psi \right] e^{\alpha q^2 t'} dt' \right) \tag{16}$$

where,

$$\Psi = \frac{P_n(a)}{k_1} f_1 - \frac{P_n(-a)}{k_2} f_2 + \frac{m\pi}{b} [(-1)^{m+1} f_4 + f_3]$$

Now, applying inversion of Fourier Cosine transform, Finite Fourier sine transform and finite Marchi-Fasulo transform to the equation (1), one obtains the expression for temperature distribution as

$$T(x, y, z, t) = \left(\frac{4\eta}{b\pi} \right) \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} \left(\frac{P_n(x)}{\lambda_n} \right) \sin py \Lambda(z) \tag{17}$$

where,

$$\Lambda(z) = \int_0^{\infty} B(t) \cos(\eta z) dz,$$

$$B(t) = e^{-\alpha q^2 t} \left(\bar{f}^* + \int_0^t \left[\frac{\alpha g}{k} + \Psi \right] e^{\alpha q^2 t'} dt' \right)$$

$$p = \left(\frac{m\pi}{b} \right)$$

$$q^2 = \left(1 + \lambda_n^2 + \frac{m^2 \pi^2}{b^2} \right)$$

Equation (2) is the required solution.

IV. AIRY'S STRESS FUNCTIONS

Substituting the value of temperature distribution $T(x,y,z,t)$ from (2) in equation (4) one obtains

$$U(x, y, z, t) = - \left(\frac{4\eta\pi E}{b\pi} \right) \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} \left(\frac{P_n(x)}{\lambda_n} \right) \sin py \Lambda(z) \tag{18}$$

Where

$$\Lambda(z) = \int_0^{\infty} B(t) \cos(\eta z) dz$$

V. DISPLACEMENT COMPONENTS

Substituting the values of Airy's stress function from equation (1) in the equation (1) to (3), one obtains

$$u_x = -\left(\frac{4\eta\lambda}{b\pi}\right)\Lambda(z) \int_{-a}^a \left\{ \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} \left(\frac{P_n(x)}{\lambda_n}\right) (-p^2 \sin py) - \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} -\eta^2 \left(\frac{P_n(x)}{\lambda_n}\right) \sin py \right. \\ \left. - \nu \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} \left(\frac{P_n''(x)}{\lambda_n^2}\right) \sin py + \lambda \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} \left(\frac{P_n(x)}{\lambda_n}\right) \sin py \right\} dx \tag{19}$$

$$u_y = -\left(\frac{4\eta\lambda}{b\pi}\right)\Lambda(z) \int_0^b \left\{ \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} \eta^2 \left(\frac{P_n(x)}{\lambda_n}\right) \sin py - \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} \left(\frac{P_n''(x)}{\lambda_n^2}\right) \sin py \right. \\ \left. + \nu \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} \left(\frac{P_n(x)}{\lambda_n}\right) (-p^2 \sin py) + \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} \left(\frac{P_n(x)}{\lambda_n}\right) \sin py \right\} dy \tag{20}$$

$$u_z = -\left(\frac{4\eta\lambda}{b\pi}\right)\Lambda(z) \int_0^{\infty} \left\{ \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} \left(\frac{P_n''(x)}{\lambda_n^2}\right) \sin py + \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} \left(\frac{P_n(x)}{\lambda_n}\right) (-p^2 \sin py) \right. \\ \left. + \nu \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} \eta^2 \left(\frac{P_n(x)}{\lambda_n}\right) \sin py - \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} \left(\frac{P_n(x)}{\lambda_n}\right) \sin py \right\} dz \tag{21}$$

VI. DETERMINATION OF STRESS FUNCTION

Substituting the value of Airy's stress function U(x,y,z,t) from equation (1) in the equation (12) to (14) one obtain the stress functions as,

$$\sigma_{xx} = -\left(\frac{4\eta\lambda E}{b\pi}\right)\Lambda(z) \sum_{n=1}^{\infty} \left(\frac{P_n(x)}{\lambda_n}\right) \\ \times \left[\sum_{m=1}^{\infty} (-p^2 \sin py) - \eta^2 \sum_{m=1}^{\infty} \sin py \right] \tag{22}$$

$$\sigma_{yy} = -\left(\frac{4\eta\lambda E}{b\pi}\right)\Lambda(z) \sum_{m=1}^{\infty} \sin py \\ \times \left[\sum_{n=1}^{\infty} \left(\frac{P_n''(x)}{\lambda_n^2}\right) - \eta^2 \sum_{n=1}^{\infty} \left(\frac{P_n(x)}{\lambda_n}\right) \right] \tag{23}$$

$$\sigma_{zz} = -\left(\frac{4\eta\lambda E}{b\pi}\right)\Lambda(z) \left[\sum_{m=1}^{\infty} \sum_{n=1}^{\infty} \left(\frac{P_n''(x)}{\lambda_n^2}\right) \sin py \right. \\ \left. + \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} \left(\frac{P_n(x)}{\lambda_n}\right) (-p^2 \sin py) \right] \tag{24}$$

VII. SPECIAL CASE

Set

$$f(x, y, z, t) = (x-a)^2 (x+a)^2 (z + e^{-z})(e^{y-t}) \tag{25}$$

$$\bar{f}(n, y, z, t) = (z + e^{-z})(e^{y-t}) \\ \times \left[\frac{a_n \cos^2(a_n a) - \cos(a_n a) \sin(a_n a)}{a_n^2} \right] \tag{26}$$

Substituting the above value in equations (2) to (3) one obtains

$$T(x, y, z, t) = \left(\frac{4\eta}{b\pi}\right) \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} \left(\frac{P_n(x)}{\lambda_n}\right) \sin py \Lambda(z) \\ \times \left[\frac{a_n \cos^2(a_n a) - \cos(a_n a) \sin(a_n a)}{a_n^2} \right] \\ (z + e^{-z})(e^{y-t}) \tag{27}$$

$$U(x, y, z, t) = -\frac{4\eta\pi E}{b\pi} \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} \left(\frac{P_n(x)}{\lambda_n}\right) \sin py \Lambda(z) \\ \times \left[\frac{a_n \cos^2(a_n a) - \cos(a_n a) \sin(a_n a)}{a_n^2} \right] \\ (z + e^{-z})(e^{y-t}) \tag{28}$$

$$u_x = -\frac{4\eta\lambda}{b\pi} \Lambda(z) \int_{-a}^a \left\{ \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} \left(\frac{P_n(x)}{\lambda_n}\right) (-p^2 \sin py) - \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} -\eta^2 \left(\frac{P_n(x)}{\lambda_n}\right) \sin py \right. \\ \left. - \nu \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} \left(\frac{P_n''(x)}{\lambda_n^2}\right) \sin py + \lambda \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} \left(\frac{P_n(x)}{\lambda_n}\right) \sin py \right\} dx \\ \times \left[\frac{a_n \cos^2(a_n a) - \cos(a_n a) \sin(a_n a)}{a_n^2} \right] \\ (z + e^{-z})(e^{y-t}) \tag{29}$$

$$u_y = -\frac{4\eta\lambda}{b\pi} \Lambda(z) \int_0^b \left\{ \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} \eta^2 \left(\frac{P_n(x)}{\lambda_n} \right) \sin py - \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} \left(\frac{P_n''(x)}{\lambda_n^2} \right) \sin py \right. \\ \left. + \nu \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} \left(\frac{P_n(x)}{\lambda_n} \right) (-p^2 \sin py) + \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} \left(\frac{P_n(x)}{\lambda_n} \right) \sin py \right\} dy \\ \times \left[\frac{a_n \cos^2(a_n a) - \cos(a_n a) \sin(a_n a)}{a_n^2} \right] (z + e^{-z})(e^{y-t}) \quad (30)$$

VIII. NUMERICAL RESULTS

Set $a = 2, k = 0.86, b = 3, t = 1$ sec in the equations (5.7.3)- (5.7.10) to obtain

$$T(x, y, z, t) = \left(\frac{2\eta}{3\pi} \right) \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} \left(\frac{P_n(x)}{\lambda_n} \right) \sin py \Lambda(z) \\ \times \left[\frac{a_n \cos^2(2a_n) - \cos(2a_n) \sin(2a_n)}{a_n^2} \right] (z + e^{-z})(e^{y-t}) \quad (35)$$

$$u_z = -\frac{4\eta\lambda}{b\pi} \Lambda(z) \int_0^{\infty} \left\{ \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} \left(\frac{P_n''(x)}{\lambda_n^2} \right) \sin py + \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} \left(\frac{P_n(x)}{\lambda_n} \right) (-p^2 \sin py) \right. \\ \left. + \nu \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} \eta^2 \left(\frac{P_n(x)}{\lambda_n} \right) \sin py - \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} \left(\frac{P_n(x)}{\lambda_n} \right) \sin py \right\} dz \\ \times \left[\frac{a_n \cos^2(a_n a) - \cos(a_n a) \sin(a_n a)}{a_n^2} \right] (z + e^{-z})(e^{y-t}) \quad (31)$$

$$U(x, y, z, t) = -\left(\frac{2\eta\pi E}{3\pi} \right) \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} \left(\frac{P_n(x)}{\lambda_n} \right) \sin py \Lambda(z) \\ \times \left[\frac{a_n \cos^2(2a_n) - \cos(2a_n) \sin(2a_n)}{a_n^2} \right] \\ (z + e^{-z})(e^{y-t}) \quad (36)$$

$$\sigma_{xx} = -\frac{4\eta\lambda E}{b\pi} \Lambda(z) \sum_{n=1}^{\infty} \left(\frac{P_n(x)}{\lambda_n} \right) \left[\sum_{m=1}^{\infty} (-p^2 \sin py) - \eta^2 \sum_{m=1}^{\infty} \sin py \right] \\ \times \left[\frac{a_n \cos^2(a_n a) - \cos(a_n a) \sin(a_n a)}{a_n^2} \right] \\ (z + e^{-z})(e^{y-t}) \quad (32)$$

$$u_x = -\left(\frac{2\eta\lambda}{3\pi} \right) \Lambda(z) \int_{-a}^a \left\{ \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} \left(\frac{P_n(x)}{\lambda_n} \right) (-p^2 \sin py) - \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} \eta^2 \left(\frac{P_n(x)}{\lambda_n} \right) \sin py \right. \\ \left. - \nu \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} \left(\frac{P_n''(x)}{\lambda_n^2} \right) \sin py + \lambda \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} \left(\frac{P_n(x)}{\lambda_n} \right) \sin py \right\} dx \\ \times \left[\frac{a_n \cos^2(2a_n) - \cos(2a_n) \sin(2a_n)}{a_n^2} \right] \\ (z + e^{-z})(e^{y-t}) \quad (37)$$

$$\sigma_{yy} = -\frac{4\eta\lambda E}{b\pi} \Lambda(z) \sum_{m=1}^{\infty} \sin py \left[\sum_{n=1}^{\infty} \left(\frac{P_n''(x)}{\lambda_n^2} \right) - \eta^2 \sum_{n=1}^{\infty} \left(\frac{P_n(x)}{\lambda_n} \right) \right] \\ \times \left[\frac{a_n \cos^2(a_n a) - \cos(a_n a) \sin(a_n a)}{a_n^2} \right] (z + e^{-z})(e^{y-t}) \quad (33)$$

$$u_y = -\left(\frac{2\eta\lambda}{3\pi} \right) \Lambda(z) \int_0^b \left\{ \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} \eta^2 \left(\frac{P_n(x)}{\lambda_n} \right) \sin py - \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} \left(\frac{P_n''(x)}{\lambda_n^2} \right) \sin py \right. \\ \left. + \nu \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} \left(\frac{P_n(x)}{\lambda_n} \right) (-p^2 \sin py) + \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} \left(\frac{P_n(x)}{\lambda_n} \right) \sin py \right\} dy \\ \times \left[\frac{a_n \cos^2(2a_n) - \cos(2a_n) \sin(2a_n)}{a_n^2} \right] \\ (z + e^{-z})(e^{y-t}) \quad (38)$$

$$u_z = -\left(\frac{2\eta\lambda}{3\pi}\right)\Lambda(z)\int_0^\infty \left\{ \sum_{m=1}^\infty \sum_{n=1}^\infty \left(\frac{P_n''(x)}{\lambda_n^2}\right) \sin py + \sum_{m=1}^\infty \sum_{n=1}^\infty \left(\frac{P_n(x)}{\lambda_n}\right) (-p^2 \sin py) \right. \\ \left. + v \sum_{m=1}^\infty \sum_{n=1}^\infty \eta^2 \left(\frac{P_n(x)}{\lambda_n}\right) \sin py - \sum_{m=1}^\infty \sum_{n=1}^\infty \left(\frac{P_n(x)}{\lambda_n}\right) \sin py \right\} dz \\ \times \left[\frac{a_n \cos^2(2a_n) - \cos(2a_n) \sin(2a_n)}{a_n^2} \right] \\ (z + e^{-z})(e^{y-t}) \quad (39)$$

$$\sigma_{xx} = -\left(\frac{2\eta\lambda E}{3\pi}\right)\Lambda(z)\sum_{n=1}^\infty \left(\frac{P_n(x)}{\lambda_n}\right) \left[\sum_{m=1}^\infty (-p^2 \sin py) - \eta^2 \sum_{m=1}^\infty \sin py \right] \\ \times \left[\frac{a_n \cos^2(2a_n) - \cos(2a_n) \sin(2a_n)}{a_n^2} \right] \\ (z + e^{-z})(e^{y-t}) \quad (40)$$

$$\sigma_{yy} = -\left(\frac{2\eta\lambda E}{3\pi}\right)\Lambda(z)\sum_{m=1}^\infty \sin py \left[\sum_{n=1}^\infty \left(\frac{P_n''(x)}{\lambda_n^2}\right) - \eta^2 \sum_{n=1}^\infty \left(\frac{P_n(x)}{\lambda_n}\right) \right] \\ \times \left[\frac{a_n \cos^2(2a_n) - \cos(2a_n) \sin(2a_n)}{a_n^2} \right] \\ (z + e^{-z})(e^{y-t}) \quad (41)$$

$$\sigma_{zz} = -\left(\frac{2\eta\lambda E}{3\pi}\right)\Lambda(z) \left[\sum_{m=1}^\infty \sum_{n=1}^\infty \left(\frac{P_n''(x)}{\lambda_n^2}\right) \sin py + \sum_{m=1}^\infty \sum_{n=1}^\infty \left(\frac{P_n(x)}{\lambda_n}\right) (-p^2 \sin py) \right] \\ \times \left[\frac{a_n \cos^2(2a_n) - \cos(2a_n) \sin(2a_n)}{a_n^2} \right] \\ (z + e^{-z})(e^{y-t}) \quad (42)$$

IX. MATERIAL PROPERTIES

The numerical calculations has been carried out for an Aluminum (pure) rectangular beam with the material properties as,

Density $\rho = 169 \text{ lb/ft}^3$

Specific heat = 0.208 Btu/lbOF

Thermal conductivity $K = 117 \text{ Btu/(hr. ftOF)}$

Thermal diffusivity $\alpha = 3.33 \text{ ft}^2/\text{hr.}$

Poisson ratio $\nu = 0.35$

Coefficient of linear thermal expansion $\alpha_t = 12.84 \times 10^{-6}/\text{F}$

Lame constant $\mu = 26.67$

Young's modulus of elasticity $E = 70 \text{ G Pa}$

X. DIMENSIONS

The constants associated with the numerical calculation are taken as

Length of rectangular beam $x = 4 \text{ ft}$

Breath of rectangular beam $y = 3 \text{ ft}$

Height of rectangular beam $z = 10^3 \text{ ft}$

XI. CONCLUSION

The temperature distribution, displacements and thermal stresses at any point of a semi-infinite rectangular beam have been obtained, when the boundary conditions are known with the aid of finite Marchi-Fasulo transform, finite Fourier sine and Fourier cosine transform techniques. The results are obtain in the form of infinite series.

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